

A NUMERICAL STUDY OF THE EVOLUTION OF THE REAL TIME IN QUANTUM
MECHANICS FOR GAUGE THEORY (QUARKS AND GLUONS PLASMA)

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Abstract

In this paper, we take the effective Hamiltonian operator until the sixth degree [3] and we apply the perturbation theory (that depends on creation operator D_i^\dagger and annihilation operator D_i) on remaining homogenous modes after quantization of the inhomogenous modes and we have concluded the time evolution for the average ensemble of global square operator (the color magnetic energy E_{mag}) analytically, and we calculated this time evolution with MATLAB software, we found that life quark and gluons plasma of rank 10^{-29} sec According to coupling constant values.

Keywords: Real times in non-equilibrium, phase transition to quark-gluon-plasma.

1-Introduction

Understanding the evolution of real time in problems of non-equilibrium, is very important, that is manifested for example in identification on the process of early universe formation, or in the theoretical study of particles under limited conditions, Aimed to predicting a very short transition of phase of quarks and gluons plasma. These problems are treated within the framework of the groups SU(2) and SU(3) in gauge theories [1-4].

-There are two methods used for processing problems by real time:

1- The first method: depends on Schrodinger representation in quantum mechanics, where the operators are not dependent on time [5,6].

2- The second method: depends on Heisenberg representation in quantum mechanics, where the operators are dependent on time $A_H(t)$, and is a processed problems of non- equilibrium, either depending on Green function, or Wigner way [4,7,8].

_The real time evolution of quarks and gluons plasma was studied for the pure gauge theory with the two groups SU(2) and SU(3) in [5-6]. In these studies, the perturbation theory which

depends on the creation operator \hat{D}^\dagger and annihilation operator \hat{D} , was used. The effective Hamiltonian operator expansion was taken into consideration until the fourth degree.

-In [13] the harmonic oscillator of pure gauge theory with group SU(2) was studied numerically using creation operator \hat{D}^\dagger and annihilation operator \hat{D} .

-In [14] We take the effective Hamiltonian operator until the sixth degree with the group SU(2); which represents nine anharmonic oscillators. We used the perturbation theory and two operators of creation and annihilation then calculate and draw the energy levels from E_{10} to E_{11} at which this study is still correct and stable. We conclude that the perturbation theory is not broken down when $g \geq 0.549915534$; it still correct until $g = 3.998312073$.

-in this work, the starting point is gauge theory (quarks and gluons plasma) and taking the effective Hamiltonian operator expansion until the sixth degree and we used perturbation theory which depends on creation operator \hat{D}_i^\dagger and annihilation operator \hat{D}_i .

2-RESEARCH METHODOLOGY:

“ Introduction to our research in words ”

According to [3], the Hamilton operator of gauge theory with group SU(2) can be described in loop (L^3)

$$\begin{aligned} \hat{H}_{eff} = & \frac{1}{2} \left(\frac{1}{g^2(L)} + \alpha_1 + n_1 f_1 \right) \hat{P}_i \hat{P}_i + (\alpha_1 + n_1 f_1) \hat{P}_i \hat{P}_i + \frac{1}{2} \left(\frac{1}{g^2(L)} + \alpha_1 + n_1 f_1 \right) \hat{P}_i (\hat{E}_i \hat{E}_i) \\ & + (\alpha_2 + n_2 f_2) \hat{P}_i \hat{P}_i \hat{P}_i \hat{P}_i + (\alpha_2 + n_2 f_2) \hat{P}_i \hat{P}_i \hat{P}_i \hat{P}_i + (\alpha_2 + n_2 f_2) \sum_i (\hat{P}_i \hat{P}_i)^2 \\ & + (\alpha_3 + n_3 f_3) \sum_i (\hat{P}_i \hat{P}_i)^2 \hat{P}_i \hat{P}_i + (\alpha_3 + n_3 f_3) \hat{P}_i \hat{P}_i \hat{P}_i \hat{P}_i \hat{P}_i \hat{P}_i \\ & + (\alpha_4 + n_4 f_4) \hat{P}_i \hat{P}_i (\hat{E}_i \hat{E}_i) \hat{P}_i \hat{P}_i + (\alpha_4 + n_4 f_4) \hat{P}_i (\hat{E}_i \hat{E}_i) \hat{P}_i \hat{P}_i \\ & + (\alpha_{20} + n_{20} f_{20}) (\hat{P}_i \hat{P}_i \hat{P}_i)^2 + g(\hat{P}_i)^2 \end{aligned} \quad (1)$$

$n_f = 3$ Number of quarks flavors for group SU(2).

where $i, j=1,2,3$ the guide of local coordinates,

$a, b=1,2,3$ are the evidences of group SU(2) generators, -

$\alpha_{1,2,3}$ are constants resulted from quantization of inhomogeneous modes of gluons field by the method of paths integration (color magnetic energy).

$f_{1,2,3}$ are constants resulted from quantization of inhomogeneous modes of quarks field by the method of paths integration (color magnetic energy).

α_0 is a constant resulted from quantization of inhomogeneous time derivative modes of gluons field by the method of paths integration (color electric energy).

f_0 is a constant resulted from quantization of inhomogeneous time derivative modes of quarks field by the method of paths integration (color electric energy) and it has the following values [3]:

$$\begin{aligned}
 \alpha_1 &= 0.001910428, \alpha_2 = -0.80164661, \alpha_3 = 0.0072714388 \\
 \alpha_4 &= 0.000626204228, \alpha_5 = -0.000748827, \alpha_6 = 0.00004867666 \\
 \alpha_7 &= -0.00008217228, \alpha_8 = -0.0012428881, \alpha_9 = -0.00011189268 \\
 \alpha_{10} &= -0.00021478176, \alpha_{11} = -0.0012778632 \quad (2.a) \\
 f_1 &= -0.00006186422, f_2 = 0.042844024, f_3 = -0.0004428844 \\
 f_4 &= 0.000788842888, f_5 = -0.001888048, f_6 = 0.0000027818812 \\
 f_7 &= -0.000028187826, f_8 = 0.00018884834, f_9 = -0.0000060887872 \\
 f_{10} &= -0.000064818046, f_{11} = 0.000064848472 \quad (2.b)
 \end{aligned}$$

Note that some values constants α_m (note that m take values from 0 until 10) is different than it is in the case of pure gauge because of the mutual effect between quarks and gluons. F_{ij}^a are tensors of the magnetic field intensity represented as [3]:

$$F_{ij}^a = \epsilon^{ijk} B_j^a \quad (3)$$

B_j^a is the operator of homogeneous magnetic field, and \hat{M}_j^a is the operator of momentum.

$$\epsilon^{ijk} = \begin{cases} 1 & \text{At direct replacement} \\ 0 & \text{When two evidences are equal} \\ -1 & \text{At indirect replacement} \end{cases}$$

$O(B^a)$ indicates that the limits of degree greater than B^a are neglected.

$g^a(L)$ is a coupling constant represented in the following relation [3]:

$$g^a(L) = -g_1 \ln(L\Lambda_{ms}) + \frac{b_2 \ln[-2 \ln(L\Lambda_{ms})]}{2b_2^2} \quad (4)$$

$$b_1 = \frac{1}{(4\pi)^2} \left(\frac{11}{3} N - \frac{2}{3} n_f \right)$$

$$b_2 = \frac{1}{(4\pi)^2} \left(-\frac{84}{3} N^2 + \frac{10}{3} N n_f + (n_f - 1) n_f / 3 \right)$$

$\Lambda_{ms} = 74.1705 \text{ MeV}$ represents an identified constant by minimum subtraction of dimension organization.

$N=2$ Number of dimensions group $SU(2)$, L is the loop length in all spatial directions.

According to this method of Hamilton operator H_{eff} , the study of gauge theory with group $SU(2)$ becomes a form of quantum mechanics with group $SU(2)$, This mean that the study of infinite number of particles and freedom degrees (quarks and gluons plasma), has been physically transformed to a study of three global particles. naemly, to confine the study to nine harmonic oscillators. Nine freedom degrees and particularly nine anhamonic oscillators. , the creation operator can be identified as the following [5,6]:

$$B_j^a = \sqrt{\frac{\hbar}{2\hbar}} B_j^a - \frac{i}{\sqrt{2\hbar}} \frac{\hat{M}_j^a}{\sqrt{\hbar}} \quad (5)$$

and annihilation operator is defined as:

$$B_j^a = \sqrt{\frac{\hbar}{2\hbar}} B_j^a + \frac{i}{\sqrt{2\hbar}} \frac{\hat{M}_j^a}{\sqrt{\hbar}} \quad (6)$$

$\hbar = 1$ (Plank constant) in a system of natural units .

Where:

$$g_1 = \left(\frac{1}{g^a(L)} + \alpha_1 + \alpha_2 f_1 \right)^{-1}, g_2 = 2(\alpha_1 + \alpha_2 f_1)$$

In this case, we find that:

$$[\mathcal{B}_x, \mathcal{B}_y]_- = a_1 \mathcal{B}_z$$

$$[\mathcal{B}_x, \mathcal{B}_z]_- = [\mathcal{B}_y, \mathcal{B}_z]_- = 0$$

$\mathcal{B}_x, \mathcal{B}_y, \mathcal{B}_z$ are Kroanker symbols of spatial coordinates and evidences of generating group SU(2). These Kroanker constants are respectively, defined as:

$$a_1 = \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}$$

$$a_2 = \begin{pmatrix} 1 & a+b \\ 0 & a-b \end{pmatrix}$$

The result of adding the equation (6) to (7) is the operator of homogeneous magnetic field:

$$\mathcal{H}_1 = \sqrt{\frac{\hbar}{2}} \left(\frac{\mathcal{B}_x + \mathcal{B}_y}{\sqrt{a_1}} \right) \quad (8)$$

The result of subtracting the equation (6) out of (7) is the operator of momentum given as:

$$\mathcal{H}_2 = i \sqrt{\frac{\hbar}{2}} \left(\frac{\mathcal{B}_x - \mathcal{B}_y}{\sqrt{a_2}} \right) \quad (9)$$

3- RESULTS AND DISCUSSION:

$\mathcal{H}_{eff}^{(0)}$ is the harmonic part of the operator \mathcal{H}_{eff} :-

$$\mathcal{H}_{eff}^{(0)} = \sum_{n_1=0}^{\infty} \sum_{n_2=0}^{\infty} \left[\frac{1}{2} \left(\frac{1}{\mathcal{Q}_x \mathcal{Q}_y} + a_2 + n_2 \mathcal{E}_z \right) \mathcal{H}_1 \mathcal{H}_2 + (a_2 + n_2 \mathcal{E}_z) \mathcal{H}_1 \mathcal{H}_2 \right]$$

$$\mathcal{H}_{eff}^{(0)} = \sum_{n_1=0}^{\infty} \sum_{n_2=0}^{\infty} \left[\frac{1}{2} a_2 \mathcal{H}_1 \mathcal{H}_2 + \frac{1}{2} a_2 \mathcal{H}_1 \mathcal{H}_2 \right] \quad (10)$$

We calculated Hamiltonian operator for harmonic oscillator:

$$\mathcal{H}_{eff}^{(0)} = \hbar \omega_1 \mathcal{B}_z \left(\mathcal{H} + \frac{\mathcal{H}}{2} \right) \quad (10)$$

Following [5,6], using the equations

$$\mathcal{B}_x^+ |n_1, n_2\rangle = \sqrt{n_1} |n_1-1, n_2\rangle \quad (11)$$

$$\mathcal{B}_x^- |n_1, n_2\rangle = \sqrt{n_1+1} |n_1+1, n_2\rangle \quad (12)$$

$$\mathcal{B}_y^+ |n_1, n_2\rangle = n_2 |n_1, n_2-1\rangle \quad (13)$$

$$\mathcal{B}_y^- |n_1, n_2\rangle = 0, \mathcal{B}_z^+ |n_1, n_2\rangle = 0 \quad (14)$$

- We calculated Hamiltonian operator matrix until sixth degree $\mathcal{H}_{n_1, n_2}^{(0)}$:

$$\mathcal{H}_{n_1, n_2}^{(0)} = \langle n_1 | \mathcal{H} | n_2 \rangle$$

$$= \hbar \omega_1 \mathcal{B}_z \left(\sum_{n_1=0}^{\infty} \sum_{n_2=0}^{\infty} m_1^2 \mathcal{B}_x^+ \mathcal{B}_x^- + \frac{\mathcal{H}}{2} \right) + (a_2 + n_2 \mathcal{E}_z) \sum_{n_1=0}^{\infty} \sum_{n_2=0}^{\infty} \sum_{n_3=0}^{\infty} \left[\frac{\hbar \omega_1}{2} \right]$$

$$\left(\sqrt{m_1+1} \sqrt{m_1+2} \sqrt{m_1+3} \sqrt{m_1+4} \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- \mathcal{B}_z^+ \mathcal{B}_z^- \right)$$

$$+ m_1^2 \sqrt{m_1+1} \sqrt{m_1+2} \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- + (m_1+1) \sqrt{m_1+1} \sqrt{m_1+2} \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- +$$

$$+ m_1^2 (m_1-1) \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- + \sqrt{m_1+1} \sqrt{(m_1+2)^2} \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- +$$

$$+ (m_1)^2 \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- + m_1^2 (m_1+1) \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- + \sqrt{m_1} \sqrt{m_1-1} (m_1-2) \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- +$$

$$+ \sqrt{m_1+1} \sqrt{m_1+2} (m_1+2) \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- + m_1^2 (m_1+2) \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- +$$

$$+ (m_1+1)^2 \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- + \sqrt{m_1} \sqrt{(m_1-1)^2} \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- +$$

$$+ (m_1+1) (m_1+2) \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- + \sqrt{(m_1)^2} \sqrt{m_1-1} \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- +$$

$$+ (m_1+1) \sqrt{m_1} \sqrt{m_1-1} \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- + \sqrt{m_1} \sqrt{m_1-1} \sqrt{m_1-2} \sqrt{m_1-3} \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- +$$

$$+ (a_2 + n_2 \mathcal{E}_z) \sum_{n_1=0}^{\infty} \sum_{n_2=0}^{\infty} \sum_{n_3=0}^{\infty} \left[\frac{\hbar \omega_1}{2} \left(\sqrt{m_1+1} \sqrt{m_1+2} \sqrt{m_1+3} \sqrt{m_1+4} \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- \right) \right]$$

$$+ m_1^2 \sqrt{m_1+1} \sqrt{m_1+2} \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- + (m_1+1) \sqrt{m_1+1} \sqrt{m_1+2} \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- +$$

$$+ m_1^2 (m_1-1) \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- + \sqrt{m_1+1} \sqrt{(m_1+2)^2} \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- +$$

$$+ (m_1)^2 \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- + m_1^2 (m_1+1) \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- + \sqrt{m_1} \sqrt{m_1-1} (m_1-2) \mathcal{B}_y^+ \mathcal{B}_y^- \mathcal{B}_x^+ \mathcal{B}_x^- +$$

$$\begin{aligned}
 & +\sqrt{m^2+1}\sqrt{m^2+2}(m^2+2)\theta^{2k}\theta_{n,m}^k + m^2(m^2+1)\theta^{2k}\theta_{n,m}^k \\
 & + (m^2+1)^2\theta^{2k}\theta_{n,m}^k + \sqrt{m^2}\sqrt{(m^2-1)^2\theta^{2k}\theta_{n,m}^k} \\
 & + (m^2+1)(m^2+2)\theta^{2k}\theta_{n,m}^k + \sqrt{(m^2)^2\sqrt{m^2-1}\theta^{2k}\theta_{n,m}^k} \\
 & + (m^2+1)\sqrt{m^2}\sqrt{m^2-1}\theta^{2k}\theta_{n,m}^k + \sqrt{m^2}\sqrt{m^2-1}\sqrt{m^2-2}\sqrt{m^2-2}\theta^{2k}\theta_{n,m}^k \\
 & + (a_n + n_r f_n) \\
 & \sum_{s=1}^k \sum_{t=1}^k \left[\frac{h^2}{\theta(\frac{h}{\theta})^2} (\sqrt{m^2+1}\sqrt{m^2+2}\sqrt{m^2+2}\sqrt{m^2+4}\sqrt{m^2+2}\sqrt{m^2+6}\theta_{n,m}^k) \right. \\
 & + m^2\sqrt{m^2+1}\sqrt{m^2+2}\sqrt{m^2+2}\sqrt{m^2+4}\theta_{n,m}^k \\
 & + \sqrt{(m^2+1)^2\sqrt{m^2+2}\sqrt{m^2+2}\sqrt{m^2+4}\theta_{n,m}^k} \\
 & + m^2(m^2-2)\sqrt{m^2+1}\sqrt{m^2+2}\theta_{n,m}^k \\
 & + \sqrt{m^2+1}\sqrt{(m^2+2)^2\sqrt{m^2+2}\sqrt{m^2+4}\theta_{n,m}^k} \\
 & + (m^2)^2\sqrt{m^2+1}\sqrt{m^2+2}\theta_{n,m}^k + m^2\sqrt{(m^2+1)^2\sqrt{m^2+2}\theta_{n,m}^k} \\
 & + m^2(m^2-1)(m^2-2)\theta_{n,m}^k + \sqrt{m^2+1}\sqrt{m^2+2}\sqrt{(m^2+2)^2\sqrt{m^2+4}\theta_{n,m}^k} \\
 & + m^2\sqrt{(m^2+1)^2\sqrt{m^2+2}\theta_{n,m}^k} + \sqrt{(m^2+1)^2\sqrt{m^2+2}\theta_{n,m}^k} \\
 & + m^2(m^2-2)^2\theta_{n,m}^k + \sqrt{(m^2+1)^2}\sqrt{(m^2+2)^2}\theta_{n,m}^k \\
 & + (m^2)^2(m^2-1)\theta_{n,m}^k + m^2(m^2+1)(m^2-1)\theta_{n,m}^k \\
 & + \sqrt{m^2}\sqrt{m^2-1}(m^2-2)\theta_{n,m}^k \\
 & + \sqrt{m^2+1}\sqrt{m^2+2}\sqrt{m^2+2}\sqrt{(m^2+4)^2}\theta_{n,m}^k \\
 & \quad + m^2\sqrt{m^2+1}\sqrt{m^2+2}\theta_{n,m}^k + \sqrt{(m^2+1)^2}\sqrt{(m^2+2)^2}\theta_{n,m}^k \\
 & + (m^2)^2(m^2-1)\theta_{n,m}^k + \sqrt{(m^2+1)}\sqrt{(m^2+2)^2}\theta_{n,m}^k + (m^2)^2\theta_{n,m}^k \\
 & + (m^2)^2(m^2+1)\theta_{n,m}^k + \sqrt{m^2}\sqrt{m^2-1}(m^2-2)\theta_{n,m}^k \\
 & + \sqrt{m^2+1}\sqrt{(m^2+2)^2}(m^2+2)\theta_{n,m}^k + (m^2)^2(m^2+1)\theta_{n,m}^k \\
 & + m^2(m^2+1)^2\theta_{n,m}^k + \sqrt{m^2}\sqrt{(m^2-1)^2}(m^2-2)\theta_{n,m}^k \\
 & + m^2(m^2+1)(m^2+2)\theta_{n,m}^k + \sqrt{(m^2)^2\sqrt{m^2-1}(m^2-2)\theta_{n,m}^k} \\
 & \quad + (m^2+1)\sqrt{m^2}\sqrt{m^2-1}(m^2-2)\theta_{n,m}^k + \sqrt{m^2}\sqrt{m^2-1}\sqrt{m^2-2}\sqrt{m^2-2}\theta_{n,m}^k \\
 & \quad + \sqrt{m^2+1}\sqrt{m^2+2}\sqrt{m^2+2}\sqrt{m^2+4}(m^2+2)\theta_{n,m}^k \\
 & + m^2\sqrt{m^2+1}\sqrt{m^2+2}(m^2+2)\theta_{n,m}^k + \sqrt{(m^2+1)^2}\sqrt{m^2+2}(m^2+2)\theta_{n,m}^k \\
 & + m^2(m^2-1)(m^2+1)\theta_{n,m}^k + \sqrt{m^2+1}\sqrt{(m^2+2)^2}(m^2+2)\theta_{n,m}^k \\
 & + (m^2)^2(m^2+1)\theta_{n,m}^k + m^2(m^2+1)^2\theta_{n,m}^k + \sqrt{m^2}\sqrt{(m^2-1)^2}(m^2-2)\theta_{n,m}^k \\
 & + \sqrt{m^2+1}\sqrt{m^2+2}\sqrt{m^2+2}\sqrt{m^2+4}\theta_{n,m}^k + m^2(m^2+1)^2\theta_{n,m}^k \\
 & + \sqrt{m^2}\sqrt{(m^2-1)^2}\theta_{n,m}^k + (m^2+1)^2(m^2+2)\theta_{n,m}^k + \sqrt{(m^2)^2}\sqrt{(m^2-1)^2}\theta_{n,m}^k \\
 & + (m^2+1)\sqrt{m^2}\sqrt{(m^2-1)^2}\theta_{n,m}^k + \sqrt{m^2}\sqrt{m^2-1}\sqrt{m^2-2}\sqrt{(m^2-2)^2}\theta_{n,m}^k \\
 & + \sqrt{m^2+1}\sqrt{m^2+2}\sqrt{(m^2+2)^2}\theta_{n,m}^k + m^2(m^2+1)(m^2+2)\theta_{n,m}^k \\
 & + (m^2+1)^2(m^2+2)\theta_{n,m}^k + \sqrt{(m^2)^2}\sqrt{(m^2-1)^2}\theta_{n,m}^k \\
 & + (m^2+1)(m^2+2)^2\theta_{n,m}^k + \sqrt{(m^2)^2\sqrt{m^2-1}\theta_{n,m}^k} \\
 & + \sqrt{(m^2)^2\sqrt{m^2-1}(m^2+1)\theta_{n,m}^k} + \sqrt{m^2}\sqrt{m^2-1}\sqrt{(m^2-2)^2}\sqrt{m^2-2}\theta_{n,m}^k \\
 & + (m^2+1)(m^2+2)(m^2+2)\theta_{n,m}^k + \sqrt{(m^2)^2\sqrt{m^2-1}(m^2+1)\theta_{n,m}^k}
 \end{aligned}$$

$$\begin{aligned}
 & \left(m_{\mu}^2(m_{\nu}^2 + 2)\sqrt{m_{\mu}^2 m_{\nu}^2} - 2 \theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\nu}^2 \theta_{\mu\nu}^2 \right) \\
 & \left(m_{\mu}^2 \sqrt{m_{\nu}^2 m_{\rho}^2} - 1(m_{\mu}^2 + 1)\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \right) \\
 & \left(m_{\mu}^2 \sqrt{m_{\nu}^2 m_{\rho}^2} - 1\sqrt{m_{\mu}^2 m_{\nu}^2} - 1\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \right) \\
 & \left(m_{\mu}^2 + 1\right)\sqrt{m_{\nu}^2 m_{\rho}^2} + 1\sqrt{m_{\mu}^2 m_{\nu}^2} + 2\sqrt{m_{\mu}^2 m_{\rho}^2} + 2\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \\
 & \left(m_{\mu}^2 + 1\right)\sqrt{m_{\nu}^2 m_{\rho}^2} + 1\sqrt{m_{\mu}^2 m_{\nu}^2} + 2m_{\mu}^2 \theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \\
 & \left(m_{\mu}^2 + 1\right)m_{\nu}^2 \sqrt{m_{\rho}^2} + 1\sqrt{m_{\mu}^2 m_{\nu}^2} + 2\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \\
 & \left(m_{\mu}^2 + 1\right)m_{\nu}^2 m_{\rho}^2 \theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \\
 & \left(\sqrt{m_{\mu}^2 m_{\nu}^2} - 1\sqrt{m_{\mu}^2 m_{\rho}^2} + 1\sqrt{m_{\nu}^2 m_{\rho}^2} + 2\sqrt{m_{\mu}^2 m_{\nu}^2} + 1\sqrt{m_{\mu}^2 m_{\rho}^2} + 2\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \right) \\
 & \left(\sqrt{m_{\mu}^2 m_{\nu}^2} - 1\sqrt{m_{\mu}^2 m_{\rho}^2} + 1\sqrt{m_{\nu}^2 m_{\rho}^2} + 2m_{\mu}^2 \theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \right) \\
 & \left(\sqrt{m_{\mu}^2 m_{\nu}^2} - 1m_{\mu}^2 \sqrt{m_{\rho}^2} + 1\sqrt{m_{\mu}^2 m_{\nu}^2} + 2\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \right) \\
 & \left(\sqrt{m_{\mu}^2 m_{\nu}^2} - 1m_{\mu}^2 m_{\nu}^2 \theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \right) \\
 & \left(m_{\mu}^2 + 1\right)\sqrt{m_{\nu}^2 m_{\rho}^2} + 1\sqrt{m_{\mu}^2 m_{\nu}^2} + 2(m_{\mu}^2 + 1)\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \\
 & \left(m_{\mu}^2 + 1\right)\sqrt{m_{\nu}^2 m_{\rho}^2} + 1\sqrt{m_{\mu}^2 m_{\nu}^2} + 2\sqrt{m_{\mu}^2 m_{\rho}^2} - 1\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \\
 & \left(m_{\mu}^2 + 1\right)m_{\nu}^2 \sqrt{m_{\rho}^2} + 1\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \\
 & \left(m_{\mu}^2 + 1\right)m_{\nu}^2 \sqrt{m_{\rho}^2} - 1\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \\
 & \left(\sqrt{m_{\mu}^2 m_{\nu}^2} - 1\sqrt{m_{\mu}^2 m_{\rho}^2} + 1\sqrt{m_{\nu}^2 m_{\rho}^2} + 2(m_{\mu}^2 + 1)\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \right) \\
 & \left(\sqrt{m_{\mu}^2 m_{\nu}^2} - 1\sqrt{m_{\mu}^2 m_{\rho}^2} + 1\sqrt{m_{\nu}^2 m_{\rho}^2} + 2\sqrt{m_{\mu}^2 m_{\nu}^2} - 1\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \right) \\
 & \left(\sqrt{m_{\mu}^2 m_{\nu}^2} - 1m_{\mu}^2(m_{\nu}^2 + 1)\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \right) \\
 & \left(\sqrt{m_{\mu}^2 m_{\nu}^2} - 1m_{\mu}^2 \sqrt{m_{\rho}^2} - 1\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \right) \\
 & \left(m_{\mu}^2 + 1\right)(m_{\nu}^2 + 1)\sqrt{m_{\mu}^2 m_{\rho}^2} + 1\sqrt{m_{\mu}^2 m_{\nu}^2} + 2\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \\
 & \left(m_{\mu}^2 + 1\right)(m_{\nu}^2 + 1)m_{\mu}^2 \theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \\
 & \left(m_{\mu}^2 + 1\right)\sqrt{m_{\nu}^2 m_{\rho}^2} - 1\sqrt{m_{\mu}^2 m_{\nu}^2} + 1\sqrt{m_{\mu}^2 m_{\rho}^2} + 2\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \\
 & \left(m_{\mu}^2 + 1\right)\sqrt{m_{\nu}^2 m_{\rho}^2} - 1m_{\mu}^2 \theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \\
 & \left(\sqrt{m_{\mu}^2 m_{\nu}^2} - 1(m_{\mu}^2 + 2)\sqrt{m_{\nu}^2 m_{\rho}^2} + 1\sqrt{m_{\mu}^2 m_{\rho}^2} + 2\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \right) \\
 & \left(\sqrt{m_{\mu}^2 m_{\nu}^2} - 1(m_{\mu}^2 + 1)m_{\mu}^2 \theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \right) \\
 & \left(\sqrt{m_{\mu}^2 m_{\nu}^2} - 1\sqrt{m_{\mu}^2 m_{\rho}^2} - 1\sqrt{m_{\nu}^2 m_{\rho}^2} + 1\sqrt{m_{\mu}^2 m_{\nu}^2} + 2\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \right) \\
 & \left(\sqrt{m_{\mu}^2 m_{\nu}^2} - 1\sqrt{m_{\mu}^2 m_{\rho}^2} - 1m_{\mu}^2 \theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \right) \\
 & \left(m_{\mu}^2 + 1\right)(m_{\nu}^2 + 1)\sqrt{m_{\mu}^2 m_{\rho}^2} - 1\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \\
 & \left(m_{\mu}^2 + 1\right)(m_{\nu}^2 + 1)\sqrt{m_{\nu}^2 m_{\rho}^2} - 1\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \\
 & \left(m_{\mu}^2 + 1\right)\sqrt{m_{\nu}^2 m_{\rho}^2} - 1\sqrt{m_{\mu}^2 m_{\nu}^2} - 1\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \\
 & \left(\sqrt{m_{\mu}^2 m_{\nu}^2} - 1(m_{\mu}^2 + 1)(m_{\nu}^2 + 1)\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \right) \\
 & \left(\sqrt{m_{\mu}^2 m_{\nu}^2} - 1(m_{\mu}^2 + 1)\sqrt{m_{\nu}^2 m_{\rho}^2} - 1\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \right) \\
 & \left(\sqrt{m_{\mu}^2 m_{\nu}^2} - 1\sqrt{m_{\mu}^2 m_{\rho}^2} - 1(m_{\nu}^2 + 1)\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \right) \\
 & \left(\sqrt{m_{\mu}^2 m_{\nu}^2} - 1\sqrt{m_{\mu}^2 m_{\rho}^2} - 1\sqrt{m_{\nu}^2 m_{\rho}^2} - 1\theta_{\mu\nu\gamma}^2 \theta_{\mu\gamma\rho}^2 \theta_{\mu\nu}^2 \right) \tag{15}
 \end{aligned}$$

The equation (15) represents Hamiltonian operator matrix until sixth degree in case gauge theory (quarks and gluons plasma).

-To calculate the evolution of time for the average values for colored magnetic energy

$$\begin{aligned}
 \sum_{\alpha=1}^3 \sum_{\beta=1}^3 \mathbb{E}_{\alpha}^{\beta} \mathbb{E}_{\beta}^{\alpha} & \\
 \sum_{\alpha=1}^3 \sum_{\beta=1}^3 \mathbb{E}_{\alpha} \mathbb{E}_{\beta} & = \frac{\hbar}{2\sqrt{E_{\alpha}^{\beta}}} (\mathbb{E}_{\alpha}^{\beta} \mathbb{E}_{\beta}^{\alpha} + \mathbb{E}_{\beta}^{\alpha} \mathbb{E}_{\alpha}^{\beta} + \mathbb{E}_{\alpha} \mathbb{E}_{\beta} + \mathbb{E}_{\beta} \mathbb{E}_{\alpha})
 \end{aligned}$$

- Depending on (11), (12) we find:

$$\begin{aligned}
 \langle \mathbb{E}_{\alpha} | \mathbb{E}_{\beta} \mathbb{E}_{\gamma} | m \rangle & = \frac{\hbar}{2\sqrt{E_{\alpha}^{\beta}}} \left[\left(\sum_{\alpha=1}^3 \sum_{\beta=1}^3 \sqrt{m_{\alpha}^2 + 1}\sqrt{m_{\beta}^2 + 2} \theta_{\alpha\beta\gamma}^2 \right) + 2m_{\alpha}^2 \theta_{\alpha\beta\gamma}^2 \right. \\
 & \left. + \sqrt{m_{\alpha}^2 m_{\beta}^2} - 1\theta_{\alpha\beta\gamma}^2 \right] + \theta \tag{16}
 \end{aligned}$$

The equation (16) represents color magnetic energy matrix.

The colored magnetic energy matrix checks the equation:

$$i\hbar \frac{\partial \mathbb{E}_{\alpha}^{\beta} \mathbb{E}_{\beta}^{\alpha}}{\partial t} = (\mathbb{E}_{\beta}^{\alpha} \mathbb{E}_{\alpha}^{\beta}) H - H (\mathbb{E}_{\beta}^{\alpha} \mathbb{E}_{\alpha}^{\beta}) \tag{17}$$

And thus:

$$\begin{aligned}
 \frac{d}{dt} \langle |B|^2 \rangle &= \langle |B|^2 \rangle' - \langle |B|^2 \rangle'' \\
 &= \sum_{n=0}^{\infty} \left[\langle |B|^2 \rangle'_{n+1} - \langle |B|^2 \rangle'_{n-1} \right] \\
 &= \sum_{n=0}^{\infty} \left[\langle |B|^2 \rangle'_{n+1} - \langle |B|^2 \rangle'_{n-1} \right] \\
 \frac{d}{dt} \langle |B|^2 \rangle &= \sum_{n=0}^{\infty} \left[\langle |B|^2 \rangle'_{n+1} - \langle |B|^2 \rangle'_{n-1} \right] \quad (18)
 \end{aligned}$$

We can calculate the evolution of time for the average values in equation (18) numerically by MATLAB program after calculating the numerical values of matrix (15),(16) by Fortran(77) language program.

We have got [10]: $1\text{GeV}^{-1}=6.58 \times 10^{-25}$ sec

We draw graphs which represent the evolution of time of magnetic energy BB for levels $n=0,1,2$ for different values for coupling constant g we find at what time the phase change occurs and thus we get the life time of the quarks and gluons plasma:

We conclude from these forms that for small and large sizes for $L=2.656$ fm which corresponds them $g=0.513652004$ to $L=1.86$ fm which corresponds them $g=2.597495692$ is being life of quarks and gluons plasma of rank 10^{-29} sec, Because phase change occurs at this time after this time we move on to phase of the Hadrons.

4 CONCLUSIONS AND RECOMMENDATIONS:

This study is the first in quantum mechanics which takes a numerical study of the evolution of the real time for gauge theory with group SU(2) i.e. nine freedom degrees, we recommend a numerical study of the evolution of the real time for gauge theory with group SU(3) i.e. twenty four freedom degree.

REFERENCES:

1. KOLLER,J.;VAN BAAL,P.-Arigorousnonperturbative result for the glueball mass and electric flux energy in a finite volume Nncl.phys.B North-Holland. Vol.273, N^o.2,1986,387-412.
2. KRIPFGANZ , J. and MICHAEL, C.- Fermionic Contributions to The Glueball Spectrum In a Small Volume Phys. Lett.B North-Holland vol. 209, N^o. 1, 1988. 77-79.
3. Van Baal ,P, 1988 the small volume expansion of gauge theories coupled to Massless fermions. Nuclear Physics B-North-Holland Vol 307,274-290.
4. AL-CHATOURI,S.-Untersushungenzumrealzeit-verhlttenQuantenfeldtheoritischemodelle Dissertation, Leipzig uni.-1991 -, 101p.
5. Dr.AL-chatouri,salman-Evolutionof Real Times in the Problems of Non-equilibrium for Pure Gauge Theory with Group SU(2),Depending on the Creation and Annihilation Operators TishreenUniversity Journal-vol.(30)No.(1) 2008,23-45.
6. Dr.AL-chatouri,salman-Evolutionof Real Times in the Problems of Non-equilibrium for Pure Gauge Theory with Group SU(3),Depending on the Creation and Annihilation Operators. Tishreen University Journal-vol.(30)No.(3) 2008,41-61.

7. Dr.AL- chatouri,salman;AL-khassi,Silva; An analytical study of the Evolution of the Real time in statistical Quantum mechanics of the Pure Gauge theory(Gluons without Quarks) with Potential expansion until the sixth degree. Tishreen University Journal-vol.(36)No.(6) 2014,131-145.
8. Dr.AL- chatouri, salman; AL-khassi,Silva; An analytical study of the Evolution of the Real time in statistical Quantum mechanics of the Gauge theory(Quarks and Gluons) with Potential expansion until the sixth degree. Tishreen University Journal-vol.(37)No.(1) 2015,183-206.
9. PATHRIA, R.K.-Statistical Mechanics,Great Britain by BPC Wheatons Ltd, Exeter, 1995,529.
10. www.phys.ufl.edu/~korytov/.../note-01-NaturalUnits-SMsummary.pdf.
11. www.livescience.com/22320-quark-gluon-plasma-big-bang-conditions.html.
12. www.nobelprize.org/noble-prizes/physics/laureates/2004/..
13. Dr.AL- chatouri,salman;Dr.Nizam,Nohey-Aldin; AL-khassi,Silva-Harmonic Oscillator Study of Pure Gauge Theory with SU(2) Group and Glonon Semi-Particle novelty journal-vol(5), Issue 3, ; 2018 pp:(10-21).
14. Dr.AL- chatouri, salman; Dr.Nizam,Nohey-Aldin; AL-khassi,Silva-ANHARMONIC OSCILLATOR STUDY FOR PURE GAUGE THEORY (GLUONS WITHOUT QUARKS) WITH GROUP SU(2)JOURNAL OF INTERNATIONAL ACADEMIC RESEARCH FOR MULTIDISCIPLINARY-vol(6), Issue 10, ; 2018 pp:(1-18).

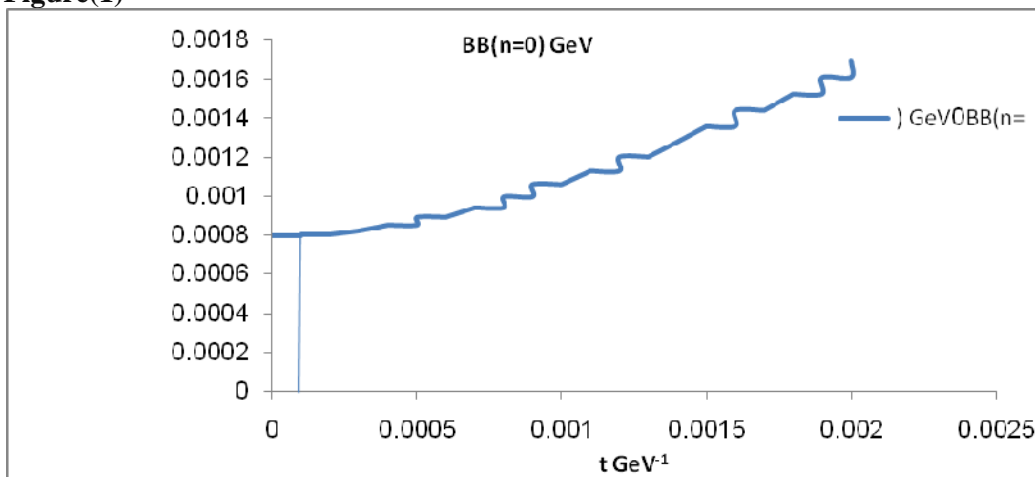
Figure legends:

Figure(1): shows the evolution of time for the average values of colored magnetic energy for ground level dependency the time for $g=0.649301461$.

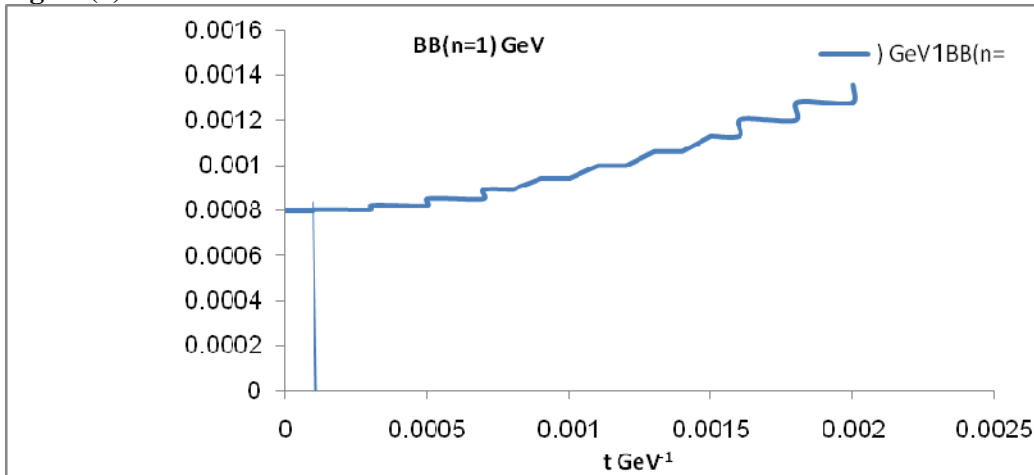
Figure(2): shows the evolution of time for the average values of colored magnetic energy for first level dependency the time for $g=0.649301461$.

Figure(3): shows the evolution of time for the average values of colored magnetic energy for second level dependency the time for $g=0.649301461$.

Figure(1)



Figure(2)



Figure(3)

